## Structural Characterisation of Organometallic Intercalates of SnSe<sub>2</sub> and ZrS<sub>2</sub> by Neutron and X-Ray Diffraction

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The synthesis of large, single crystals of ZrS<sub>2</sub> intercalated with various organometallic sandwich compounds  $\{[Co(\eta-C_5H_5)_2], [Cr(\eta-C_5H_5)_2], [Co(\eta-C_5H_4Me)_2], [Mo(\eta-C_6H_6)_2], [W(\eta-C_7H_7)(\eta-C_5H_4Me)] \text{ and } [Ti(\eta-C_8H_8)(\eta-C_5H_5)]\}$  and of SnSe<sub>2</sub> intercalated with  $[Co(\eta-C_5H_5)_2]$  are described; results from both neutron and X-ray diffraction experiments reveal that for all the intercalates the organometallic guests reside with their metal-to-ring centroid axes parallel to the layers of the host lattice.

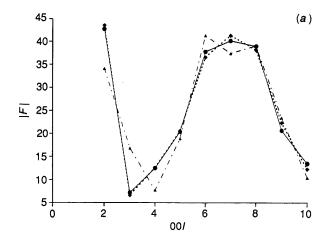
Intercalation into the interlamellar space of metal dichalcogenides modifies the properties of the host material, often quite dramatically; in the case of  $\rm SnSe_2$ , intercalation of  $\rm [Co(\eta\text{-}C_5H_5)_2]$  converts it from an anisotropic semiconductor to a relatively isotropic Type II superconductor with  $T_c=8.3~\rm K.^1$  While extensive scrutiny has been devoted to studying the physical, magnetic and electrical properties of this class of metal dichalcogenide intercalates,  $^{2-10}$  fundamental questions such as the orientational preferences of the metallocene guests in the interlamellar space still remain largely unresolved and are plagued with conflicting evidence.

Silbernagel's <sup>1</sup>H NMR study in 1975 on TaS<sub>2</sub>[Co( $\eta$ -C<sub>5</sub>H<sub>5</sub>)<sub>2</sub>]<sub>0.25</sub> concluded that the  $C_5$  axis of [Co( $\eta$ -C<sub>5</sub>H<sub>5</sub>)<sub>2</sub>] lies parallel to the layers, <sup>11</sup> an observation corroborated by Green and coworkers for the ZrS<sub>2</sub> intercalates, <sup>12</sup> where the interlayer lattice expansion was correlated with the sizes of a series of guests, varying from [Cr( $\eta$ -C<sub>5</sub>H<sub>5</sub>)<sub>2</sub>] to [Cr( $\eta$ -C<sub>5</sub>H<sub>5</sub>)( $\eta$ -C<sub>7</sub>H<sub>7</sub>)]. However, a detailed deuterium NMR study by Heyes in 1987 suggested that while the [Co( $\eta$ -C<sub>5</sub>H<sub>5</sub>)<sub>2</sub>] in TaS<sub>2</sub> exists as a

mixture of two orientations<sup>13</sup> at room temperature, the orientation with the  $C_5$  axis perpendicular to the layers is exclusively adopted at temperatures below 230 K. In contrast,  $SnS_2[Co(\eta-C_5D_5)_2]_{0.3}$  was shown, both by single crystal <sup>2</sup>H NMR studies<sup>14</sup> as well as X-ray and neutron diffraction experiments,<sup>15</sup> to consist solely of an orientation of  $[Co(\eta-C_5D_5)_2]$  with the  $C_5$  axis parallel to the layers.

We report here the synthesis of large single crystals (typically  $2 \times 2 \times 0.2$  mm) of a series of  $ZrS_2$  and  $SnSe_2$  intercalates, using both symmetrical and unsymmetrical organometallic sandwich complexes, and the application of X-ray and neutron diffraction techniques to study the guest orientation with respect to the layers.

The metal chalcogenide hosts were prepared from the elements. Elemental Zr and S (1:2, with a 1% molar excess of S) were heated to 900 °C for 1 week in evacuated, sealed silica quartz ampoules. The purple ZrS<sub>2</sub> with a reddish tinge (ZrS<sub>3</sub>) was ground under a nitrogen atmosphere and reannealed at 900 °C for another week to give a free flowing, purple-black



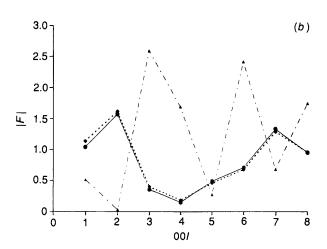


Fig. 1 Comparison of structure factors obtained from  $ZrS_2[Cr(\eta-C_5H_5)_2]_{0.20}$  ( $F_{obs}$ ) vs. structure factors calculated from models in which the  $[Cr(\eta-C_5H_5)_2]$  adopts orientations where the  $C_5$  axis lies either exclusively parallel ( $F_{para}$ ) or perpendicular ( $F_{perp}$ ) to the host layers: (a) data from neutron diffraction experiments; (b) data from X-ray diffraction experiments. (—  $F_{obs}$ , - - - -  $F_{para}$ , ----  $F_{perp}$ )

powder. The purity of the  $ZrS_2$  was checked by powder X-ray diffraction. Small quantities of  $ZrS_3$  impurities, if present, were removed by sublimation at 900 °C in an evacuated silica ampoule. Large single crystals of  $ZrS_2$  were grown by iodine vapour-phase transport of the  $ZrS_2$  powder, in a linear temperature gradient of 900 to 800 °C.  $SnSe_2$  was prepared in a similar manner but at lower reaction and growth temperatures (560 and 510 °C respectively).

Large single crystals of the intercalate were obtained by heating the crystals ( $ca. 2 \times 2 \times 0.2$  mm) in a sealed Young's ampoule containing a toluene solution of the appropriate organometallic component at 120 °C for one (ZrS<sub>2</sub> intercalates) or two weeks (SnSe<sub>2</sub> intercalates). After intercalation, the plate-like crystals, which had lost their metallic lustre and also expanded visibly along the c-axis, were washed several times with toluene and stored under nitrogen.

X-Ray diffraction patterns of aligned crystals of the intercalates were recorded in Bragg-Brantano geometry so that only 00*l* reflections were collected. Integrated intensities for the 00*l* reflections were extracted after Lorentz, polarization and absorption corrections. Room-temperature neutron diffraction spectra of aligned crystals of both ZrS<sub>2</sub> and SnSe<sub>2</sub> intercalates were acquired on the liquid and amorphous diffractometer (LAD) at the Rutherford Appleton Labora-

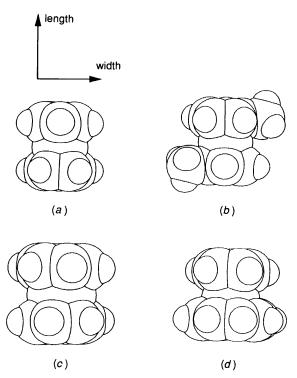


Fig. 2 Space-filling models of various guest molecules showing van der Waals diameters of the constituent atoms: (a)  $[Co(\eta-C_5H_5)_2]$ ; (b)  $[Co(\eta-C_5H_4Me)]$ ; (c)  $[Mo(\eta-C_6H_6)_2]$ ; (d)  $[Ti(\eta-C_5H_5)(\eta-C_8H_8)]$ 

tory, Didcot. The 00*l* intensities derived from both the X-ray and neutron experiments were used to perform subsequent least-squares refinements of the guest orientational parameters.

The c axis dimensions of the intercalates (obtained from least-squares analyses of the X-ray data) are tabulated in Table 1 with the c axis expansion ( $\Delta c$ ), stoichiometry and dimensions of the guest molecule. These figures compare favourably with those reported previously. 9,12,16,17

Neutron data were only obtained for the symmetrical metallocene guests  $[Co(\eta-C_5D_5)_2]$  and  $[Cr(\eta-C_5D_5)_2]$  owing to the limited access to neutron beam time and difficulty in obtaining well deuteriated samples of the other organometallic guests. However, X-ray diffraction, which is more readily accessible, was used to study a series of intercalates with guests of varying sizes and geometries, ranging from the almost spherically symmetrical  $[Co(\eta-C_5H_5)_2]$  to the larger and truncated cone-shaped  $[Ti(\eta-C_5H_5)(\eta-C_8H_8)]$  (Fig. 2).

The results of both neutron and X-ray refinements are shown in Table 2. Comparison of the observed and calculated scattering factors, as shown in Fig. 1 for the  $[Cr(\eta-C_5H_5)_2]$ intercalate of ZrS<sub>2</sub>, demonstrates the excellent level of agreement between our experimental data and that determined from a model in which all the cobaltocene molecules were oriented parallel to the host layers. The standard deviations on the refined X-ray and neutron data indicate that occupancy refinement tends to be more accurate using neutron data, while the relative S coordinates along the c axis are more precisely determined by using X-ray data refinement. This difference can be rationalized by examining the relative contributions of the guest molecule towards the overall X-ray or neutron diffraction intensity. In neutron diffraction, the guest molecules dominate the overall scattering, making the technique a much more sensitive probe of guest orientation. In contrast, X-rays are predominantly scattered by the host (ZrS<sub>2</sub> or SnSe<sub>2</sub>) and are therefore a more sensitive indicator of the changes in the S coordinates. However, while X-ray diffraction is a less discriminating probe of guest structure than neutron diffraction, comparison of

Table 1 Comparison of lattice expansion along the c axis for  $SnSe_2$  and  $ZrS_2$  intercalates with the dimensions of guest species; length and width conventions are shown in Fig. 2

Compound	c Spacinga/Å	$\Delta c/{\rm \AA}^b$	Length/Å <sup>c</sup>	Width/Å <sup>c</sup>
SnSe[Co(C <sub>5</sub> H <sub>5</sub> ) <sub>2</sub> ] <sub>0.30</sub>	11.84(1)	5.70	6.96	6.76
$ZrS_2[Co(C_5H_5)_2]_{0.25}$	11.23(1)	5.41	6.96	6.76
$ZrS_2[Cr(C_5H_5)_2]_{0.20}$	11.26(1)	5.43	$6.96^{d}$	$6.76^{d}$
$ZrS_2[Co(C_5H_4Me)_2]_{0.20}$	11.27(1)	5.44	$6.96^{d}$	7.38
$ZrS_2[W(C_7H_7)(C_5H_4Me)]_{0.20}$	11.86(3)	6.03	_	_
$ZrS_2[Mo(C_6H_6)_2]_{0.20}$	11.71(1)	5.88	7.08	7.23
$ZrS_2[Ti(C_8H_8)(C_5H_5]_{0.20}$	12.20(1)	6.37	6.97	8.15

<sup>&</sup>lt;sup>a</sup> Obtained by least-squares refinement of experimental data. <sup>b</sup> Using  $c(ZrS_2) = 5.83$  Å,  $c(SnSe_2) = 6.14$  Å. <sup>c</sup> Dimensions of unionised guests (see Fig. 2). <sup>d</sup> Based on  $[Co(\eta - C_5H_5)_2]$  dimensions.

**Table 2** Summary of occupancy refinements for both X-ray and neutron diffraction data; the R-factor is defined as  $[\Sigma(I_{\text{obs}} - I_{\text{calc}})/(\Sigma I_{\text{obs}})] \times 100$ , where  $I_{\text{obs}}$  and  $I_{\text{calc}}$  are the observed and calculated intensities respectively

Compound $(MX_2\{G\})$	X-Ray					Neutron				
	S1 coord $(z/c)$	Parallel <sup>a</sup> occupancy	Perpend.a occupancy	% parallel occupancy	R- factor (%)	\$1 coord (z/c)	Parallel <sup>a</sup> occupancy	Perpend.a occupancy	% parallel occupancy	R- factor (%)
$SnSe_2[Co(C_5H_5)_2]_{0.30}$	0.134(1)	0.29(2)	0.01(2)	96.7	1.79	0.126(4)	0.291(7)	0.009(7)	97.0	8.15
$ZrS_2[Co(C_5H_5)_2]_{0.25}$	0.136(1)	0.24(1)	0.01(1)	96.0	1.42	0.126(4)	0.242(5)	0.008(5)	96.8	5.63
$ZrS_2[Co(C_5H_5)_2]_{0.20}$	0.137(1)	0.18(1)	0.02(1)	90.0	1.28	0.129(4)	0.202(8)	-0.002(8)	100.0	5.05
$ZrS_2[Co(C_5H_4Me)_2]_{0.20}$	0.135(1)	0.17(2)	0.03(2)	85.0	3.39					
$ZrS_2[W(C_7H_7)(C_5H_4Me)_{0.20}]$	0.133(2)	0.16(5)	0.04(5)	80.0	3.51					
$ZrS_2[Mo(C_6H_6)_2]_{0.20}$	0.122(2)	0.19(2)	0.01(2)	95.0	2.70					
$ZrS_2[Ti(C_8H_8)(C_5H_5)]_{0.20}$	0.122(1)	0.19(1)	0.01(1)	95.0	1.06					

<sup>&</sup>lt;sup>a</sup> Parallel and perpendicular refer to orientations in which the metal-ring centroid axis lies parallel or perpendicular to the layers repectively.

results for  $[\text{Co}(\eta\text{-}\text{C}_5\text{H}_5)_2]$  and  $[\text{Cr}(\eta\text{-}\text{C}_5\text{H}_5)_2]$  in Table 2 shows it is still a *qualitatively* good indicator of guest orientation. Both methods indicate that the organometallic guests reside with their metal-ring centroid axes parallel to the host layers. This finding is supported by data in Table 1 which show that the *c* axis expansion increases as a function of the width of the guest, while the length remains essentially constant.

Taken together, the data in Tables 1 and 2 allow us to make several observations about the guest orientation and packing between the layers. The ring-to-ring separations of  $[Co(\eta C_5H_5)_2]$  and  $[Ti(\eta-C_5H_5)(\eta-C_8H_8)]$  are roughly the same, as is the stoichiometry of their  $ZrS_2$  intercalates. Hence the difference in c axis expansion can only be due to a parallel orientation of guest within the host layers. In contrast,  $[Co(\eta-C_5H_5)_2]$  and  $[Cr(\eta-C_5H_5)_2]$  intercalates of ZrS<sub>2</sub> show virtually the same interlayer expansions. Since they have similar molecular dimensions, the difference in stoichiometry of the intercalates reflects the absence of guest close packing in  $ZrS_2[Cr(\eta-C_5H_5)_2]_{0.20}$ . Since  $[Co(\eta-C_5H_4Me)_2]$  is somewhat bigger than  $[Co(\eta-C_5H_5)_2]$  perpendicular to the metalring centroid axis, it might be expected to induce a significantly larger c axis expansion upon intercalation into  $ZrS_2$ . The absence of such a larger expansion suggests that the methyl substituents lie between the sulfur layers, making their effective 'width' the same as  $[Co(\eta-C_5H_5)_2]$ .

In conclusion, an array of organometallic guest molecules with varying sizes and geometry have been intercalated into single crystals of ZrS<sub>2</sub>. These crystals, along with SnSe<sub>2</sub>[Co( $\eta$ -C<sub>5</sub>H<sub>5</sub>)<sub>2</sub>]<sub>0.3</sub>, have been examined by X-ray and neutron diffraction in order to extract information as to the nature of the guest orientation between the layers. While the neutron data conclusively demonstrate that both [Co( $\eta$ -C<sub>5</sub>H<sub>5</sub>)<sub>2</sub>] and [Cr( $\eta$ -C<sub>5</sub>H<sub>5</sub>)<sub>2</sub>] exist virtually exclusively in an orientation with their principal molecular axes parallel to the layers, we have shown that X-ray diffraction is a more convenient, albeit less discriminating, probe of the guest orientation between the dichalcogenide layers.

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